

Superparamagnetic properties of α - Fe_2O_3 particles: Mössbauer spectroscopy and DC magnetic measurements

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Synopsis

Mössbauer spectroscopy with ^{57}Fe

Transmission (absorption) and Conversion Electron MS

Hyperfine interactions

Zeeman splitting and relative **intensities** of lines

Hyperfine field and short range order

(Small particles and transition to superparamagnetism)

Example: Nanoparticles of Fe_2O_3 oxide

Mössbauer spectroscopy and magnetic properties

Collaborators: A. Lančok, E. Pollert, M. Maryško

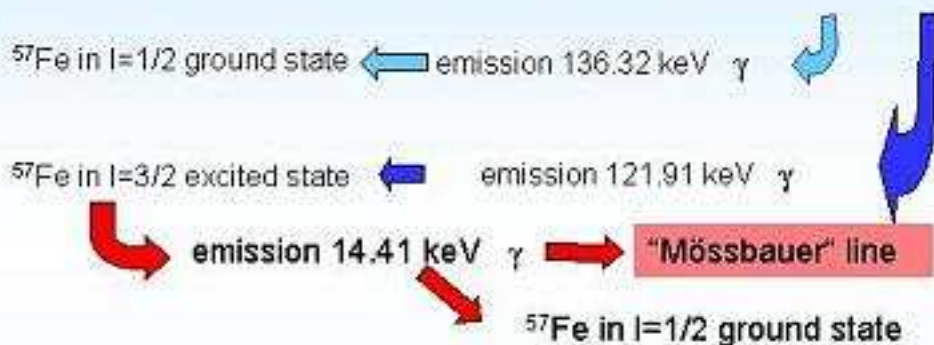
Mössbauer spectra at various temperatures

ZFC/FC magnetic moments, distribution of blocking temperatures

Mössbauer spectroscopy with ^{57}Fe

Magnetic materials containing Fe - "Mössbauer" isotope ^{57}Fe (natural abundance 2.19%)

Radioactive source ^{57}Co (electron capture with half time 270 d) \rightarrow ^{57}Fe in excited $I=7/2$ state



Absorption process



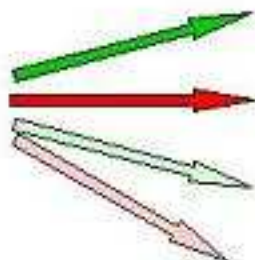
^{57}Fe in $l=3/2$ excited state



emission 14.41 keV γ



^{57}Fe in $l=1/2$ ground state



Conversion electrons K, L, M

emission 14.41 keV γ

Auger electrons KLL

X-ray $K\alpha$

Table 1. Transition from excited to ground state of the ^{57}Fe nucleus

Energies and probabilities of processes during de-excitation

type	energy	[keV]	probability	
γ emission	E_0	14.4	$1 / 1 + \alpha$	0.09
e conv K	$E_0 - B_K$	7.3	$\alpha_K / 1 + \alpha$	0.81
e conv L	$E_0 - B_L$	13.6	$\alpha_L / 1 + \alpha$	0.09
e conv M	$E_0 - B_M$	14.3	$\alpha_M / 1 + \alpha$	0.01
e Auger KLL	$B_K - 2B_L$	5.4, 5.7	$\alpha_K (1 - (FY)_K) / 1 + \alpha$	0.57
$X_{K\alpha}$	$B_K - B_L$	6.3	$\alpha_K (FY)_K / 1 + \alpha$	0.24

Hyperfine interactions

Restriction to
 ^{57}Fe

nucleus

electron shell

volume charge density

wave function inside the nucleus

electrostatic

isomer (chemical) shift

quadrupole moment

gradient of electric field

quadrupole splitting (shift)

magnetic

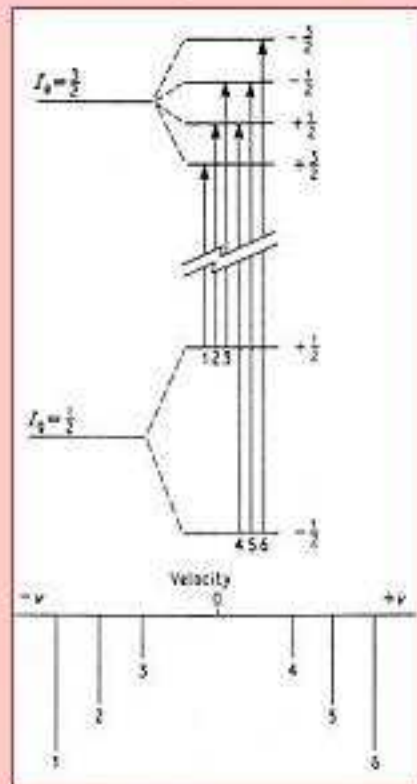
dipole moment

(effective) magnetic field inside
the nucleus (s-electrons)

nuclear Zeeman effect (splitting)

Zeeman splitting and relative intensities of lines

If the ^{57}Fe nucleus is exposed to an effective magnetic field B , its ground ($I_g = 1/2$) and excited ($I_e = 3/2$) states are split



$$E_m = -g_N \mu_N B m_z$$

g_N - nuclear Landé factor

μ_N - nuclear magneton

m_z - magnetic quantum number equal 1, 1-1, ..., -1

and 4 levels.

into 2

Transitions with Δm equal to 0 and ± 1 only allowed

spectrum consists of 6 lines

The relative intensities depend on θ

given by the Clebsch-Gordan coefficients:

$$I_1 = I_6 = 3/8 (1 + \cos^2\theta)$$

$$I_2 = I_5 = 1/2 (1 - \cos^2\theta)$$

$$I_3 = I_4 = 1/8 (1 + \cos^2\theta)$$

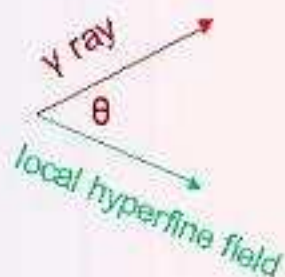
with the ratio of 2nd (5th) and 3rd (4th) lines

$$I_{2,5} / I_{3,4} = 4 \sin^2\theta / (1 + \cos^2\theta)$$



= 2 for (homogeneously) random distribution of the local moments (fields)
formally corresponding to an effective angle of 54.7°

The ratio $I_{1,6} / I_{3,4} = 3$ independently of the orientation (additional constraint for fitting).



= 0 for parallel

orientations

= 4 for perpendicular

Hyperfine field and short range order

Hyperfine field vs electronic magnetic moment:

directly



densities of the spin-up and spin-down s-electrons, $\phi^2(r) \neq 0$ in the volume of the nucleus,

indirectly



screening of s-electrons by d-electrons sensitive to the moments of d-electrons

**But generally $\mathbf{B} = f(\text{electronic magnetic moment})$
more complicated than $\mathbf{B} = \mathbf{K} \mu_e$**

B sensitive to the type and number of nearest (n.n.) and also next nearest (n.n.n.) neighbours
change of $\phi^2(r)$ caused by bonding - similar to the above case

interactions with the neighbouring magnetic moments



**B depends in a discrete way on the number of magnetic n.n.
Experimental spectrum may be decomposed to well-defined sextets**


Small particles and transition to superparamagnetism



Mössbauer Spectroscopy



Each ^{57}Fe atom with given SRO

Time average of \mathbf{B} over $10^{-7}\text{s} = 0$

Sextet collapses to doublet/singlet 

Spontaneous change of direction of 


For **superparamagnetism** decisive $K \cdot V$
Anisotropy constant K
and particle volume V

Note different characteristic time in Mössbauer Spectroscopy

and

D.C. magnetic measurements

Diffraction Methods (XR, E, N)



Coherent volume with translational symmetry

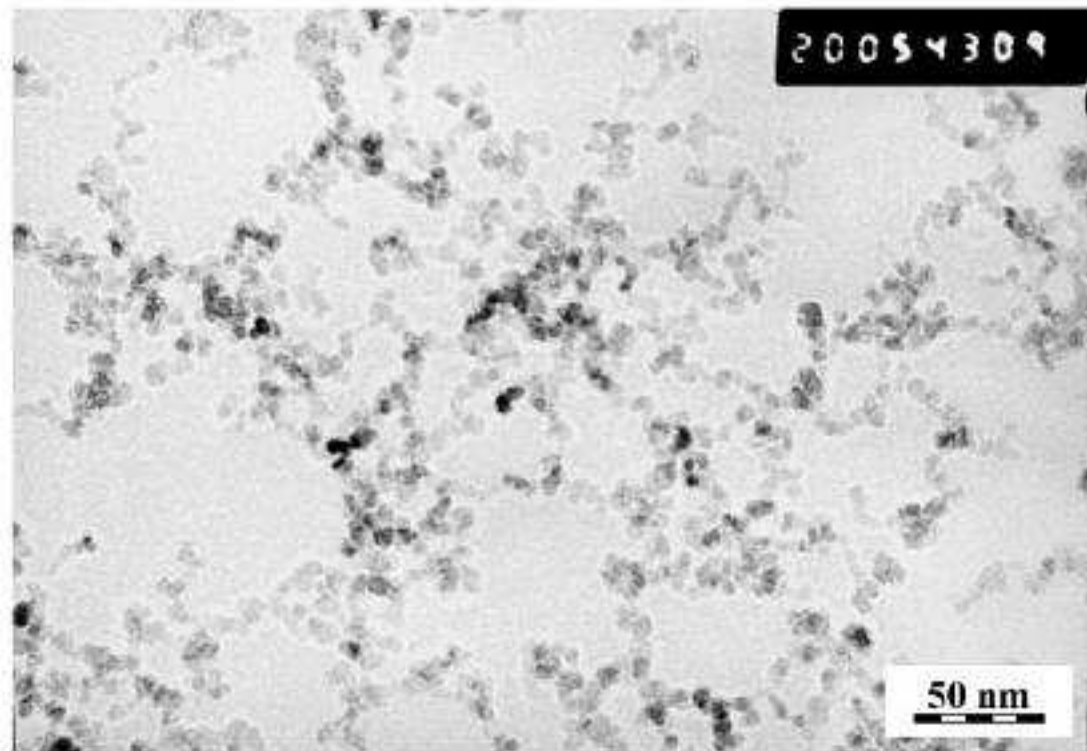
Broadening of diffraction line for small particle size

atomic moments : **paramagnetism**

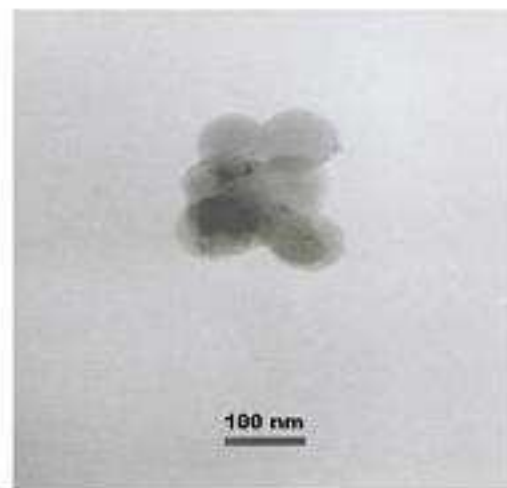
particle moments : **superparamagnetism**

vs $k_B \cdot T$ (Boltzmann constant \times Temperature)

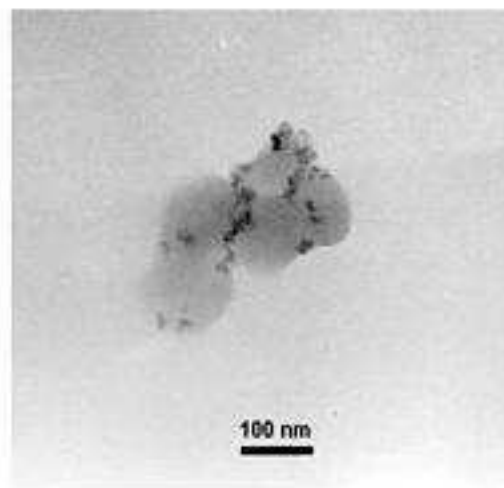
Magnetic poly(glycidyl methacrylate) microspheres containing maghemite prepared by emulsion polymerization



TEM pictures of PGMA microspheres with various concentrations of Fe_2O_3

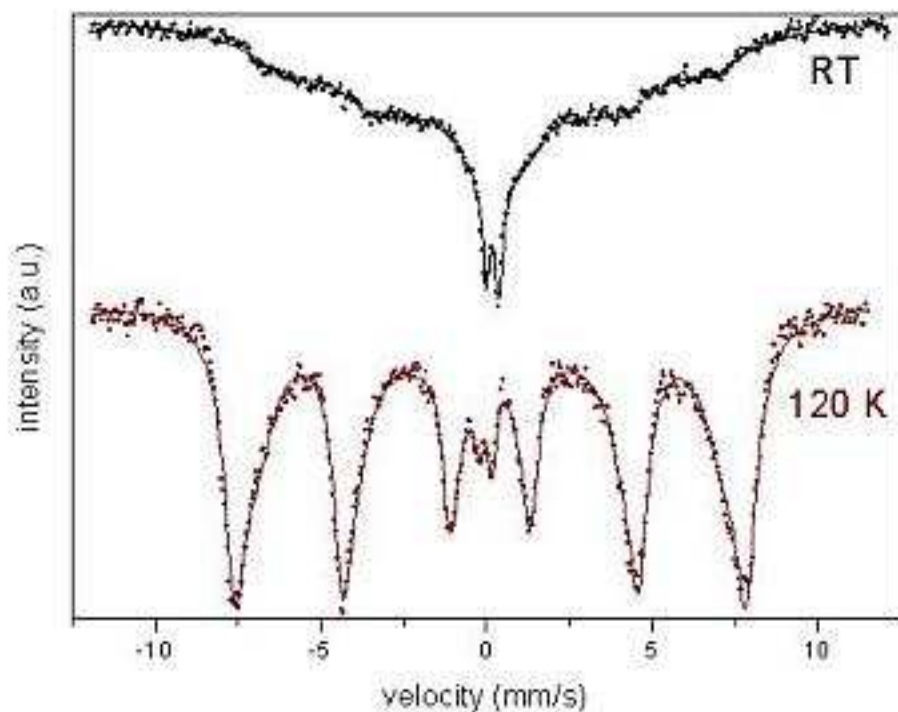


0.8 wt % of Fe_2O_3



6.2 wt % of Fe_2O_3

Moessbauer spectra at room temperature and 120 K



Parameters of the Mössbauer spectra

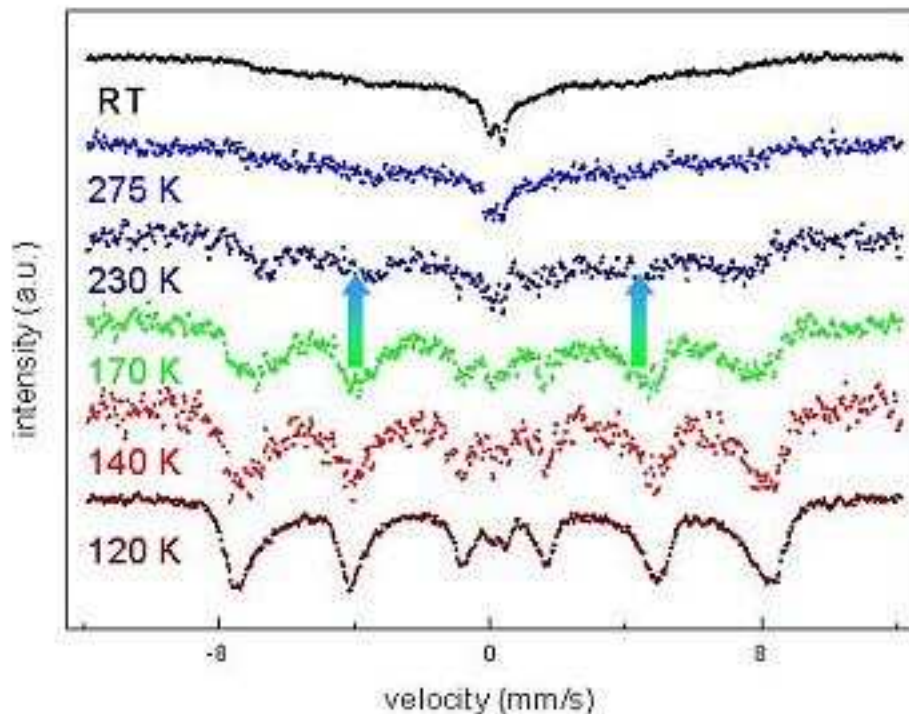
	RT			120 K		
	B_{hf}	ISO	r. a. %	B_{hf}	ISO	r.a. %
Sextet 1	40.3	0.38	14.4	47.9	0.13	38.3
Sextet 2				44.0	0.10	41.8
Broad sextet	25.9	0.29	76.0	26.4	0.34	16.7
Doublet 1		0.37	5.0		~0.0	3.2
Doublet 2		0.18	4.6			

mostly superparamagnetic
phase? Fe^{3+} !

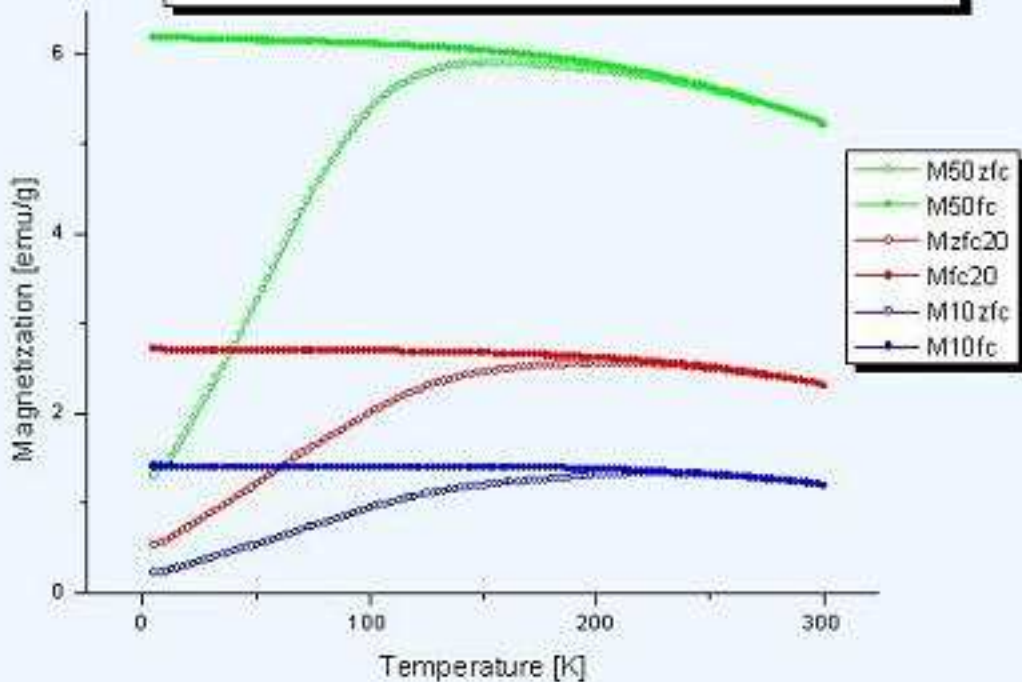
maghemite

mostly ordered

Changes of Mössbauer spectra with rising temperature



ZFC and FC magnetizations of γ -Fe₂O₃ particles



Blocking temperatures and their distribution

Even for 1 particle T_B is rather an interval – better: time of relaxation to equilibrium

time of measurement

Néel-Brown (Arrhenius-type expression) $\tau = \tau_0 \exp(\Delta E / kT)$ with $\Delta E = K \cdot V$
non-interacting, uniaxial

System of particles, way of averaging, change of ΔE by interactions

increase or decrease of T_B ?

J.J. Lu et al., JMMM 209(2000)37

claimed: distribution proportional to

$$d(M_{ZFC} - M_{FC}) / dT$$

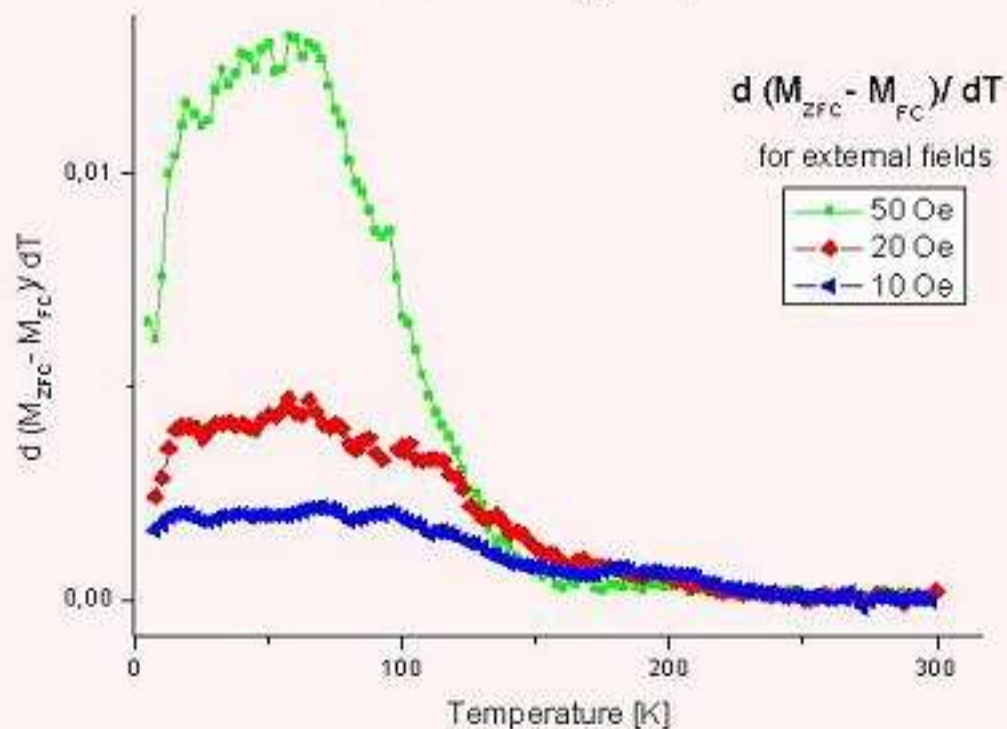
M.C. simulation for continuous distribution of particle sizes

with increasing probe field

maximum of $M_{ZFC}(T)$ shifted up


maximum of $d(\Delta M)/dT$ shifted down

Distribution of blocking temperatures



Conclusions for the Iron-oxide particles

Phase composition  not clear from XRD nor r.t. Möss. spectra

 Mössbauer spectra at 120 K indicated maghemite as main phase

Transition between ordered
and superparamagnetic state

From Mössbauer spectra
(170 – 230 K) and 10^{-7} s

From analysis of ZFC/FC magnetizations
distribution of blocking temperatures

Maxima depending on the magnetic
field intensity at (13 – 80 K)
for DC magnetizations \approx s

Moral: When speaking about superparamagnetic particles,
always state the relevant time window and temperature!

**T
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a
n
k**

Y O U !