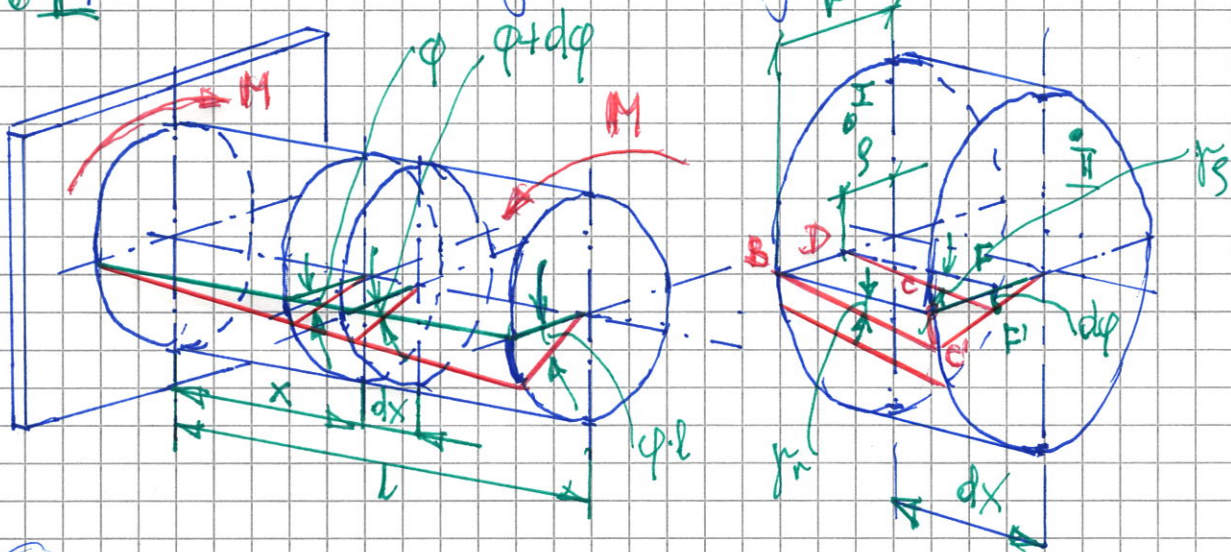


Torsion. Torsion of circular shafts Skrećenje, skrećenje krugovih prigučih

Circular shafts are commonly used for to transmit power in rotating machinery.



A) Simplifying assumptions (zapozenia)

During the deformation, the cross sections are not distorted in any manner — they remain plane, and the radius r does not change. In addition, the length l of the shaft remains constant.

Based on these observations, we make the following assumptions

- circular cross sections remain plane (do not warp) and perpendicular to the axis of the shaft,
- cross sections do not deform (there is no strain in the plane of the cross section),
- the distances between cross sections do not change (the axial normal strain is zero),

Each cross section rotates as a rigid entity about the axis of the shaft.

B) Compatibility (zgodnosć)

Because the cross sections are separated by an infinitesimal distance dx , the difference in the rotations, denoted by the $d\phi$ angle, is also infinitesimal.

As the cross sections undergo the relative rotation $d\phi$, the line BC deforms into helix BC' . Helix angle γ is the shear strain of the element.

From the geometry of figure, we obtain:

for $\theta = r$

$$\gamma_r = \frac{CC'}{BC}$$

but $BC = DF = dx$

for θ

$$\gamma = \frac{FF'}{DF}$$

$$FF' = dx \cdot \gamma \quad \text{and} \quad FF' = \theta \cdot dy$$

$$\gamma = \frac{d\theta}{dx} \cdot r$$

The corresponding shear stress is determined from Hooke's law

$$\gamma = \frac{\tau}{G}$$

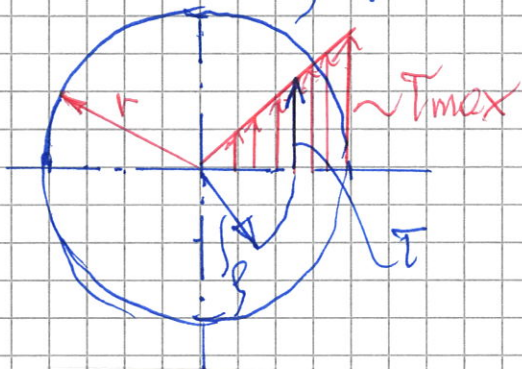
$$\tau = \gamma \cdot G = G \cdot \frac{d\theta}{dx} \cdot r$$

for selected/defined cross section

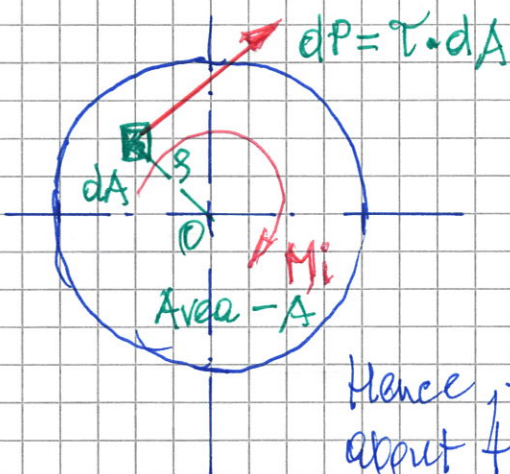
$$\frac{d\theta}{dx} = \text{const}$$

$\frac{d\theta}{dx}$ — is the angle of twist per unit length

The shear stress varies linearly with the radial distance r from the axis of the shaft.



① Equilibrium (rōnnowaga)



The shear force acting on elementary area dA is

$$dP = \tau \cdot dA = G \cdot \frac{d\varphi}{dx} \cdot \rho \cdot dA$$

directed perpendicular to the radius

Hence, the moment (torque) of the dP about the center O is

$$\rho \cdot dP = G \cdot \frac{d\varphi}{dx} \cdot \rho^2 \cdot dA$$

Summing the contributions and equating the result to the internal torque yields M_i

$$\int_A \rho dP = M_i, \text{ or}$$

$$\int_A \left(G \cdot \frac{d\varphi}{dx} \right) \cdot \rho^2 \cdot dA = M_i$$

const const.

$$G \cdot \frac{d\varphi}{dx} \int_A \rho^2 \cdot dA = M_i \quad \text{but} \quad \int_A \rho^2 dA = I_0$$

I_0 - polar moment of inertia of the cross-sectional area

$$G \cdot \frac{d\varphi}{dx} \cdot I_0 = M_i = M \quad \left(\text{but } M_i = M \text{ (external torque moment)} \right)$$

$$\frac{d\varphi}{dx} = \frac{M}{G \cdot I_0}$$

$$\varphi = \int_0^l \frac{M}{G \cdot I_0} dx \Rightarrow \frac{M \cdot l}{G \cdot I_0}$$

the angle of twist

$G I_0$ - torsional stiffness (torque)

but $\frac{dq}{dx} = \frac{T}{G \cdot r}$, and $\frac{dy}{dx} = \frac{M}{G I_0}$

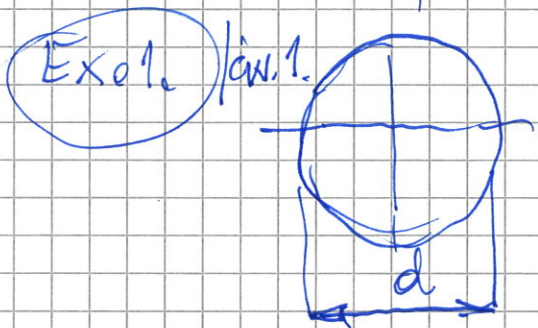
$\cancel{\tau} = \frac{dq}{dx} \cdot G \cdot r = \frac{M}{G \cdot I_0} \cdot G \cdot r = \frac{M \cdot r}{I_0}$

$\tau = \frac{M}{I_0} \cdot r$ but τ_{max} for $r = r$

$\tau_{max} = \frac{M}{I_0} \cdot r$ $\frac{\tau_0}{r} = k_{tp}$

k_{tp} - polar modulus
(section modulus for torsion)
коэффициент вытекания по сдвигу

① Calculations, examples



$\frac{M, \tau_{max}, G, l, \varphi_{max}}{d - ?}$

$I_0 = \frac{\pi d^4}{32} = \frac{\pi r^4}{2}$

$k_{tp} = \frac{\pi d^3}{16}$

$\tau_{max} \leq \frac{M}{\frac{\pi d^3}{16}}$

$\tau_{max} \leq \frac{M}{I_0} r$

$\tau_{max} \leq \frac{M}{k_{tp}}$

$\tau_{max} = \frac{16M}{\pi d^3}$

$d \geq \sqrt[3]{\frac{16M}{\pi \tau_{max}}}$

$\varphi_{max} \leq \frac{M \cdot l}{G \cdot I_0}$

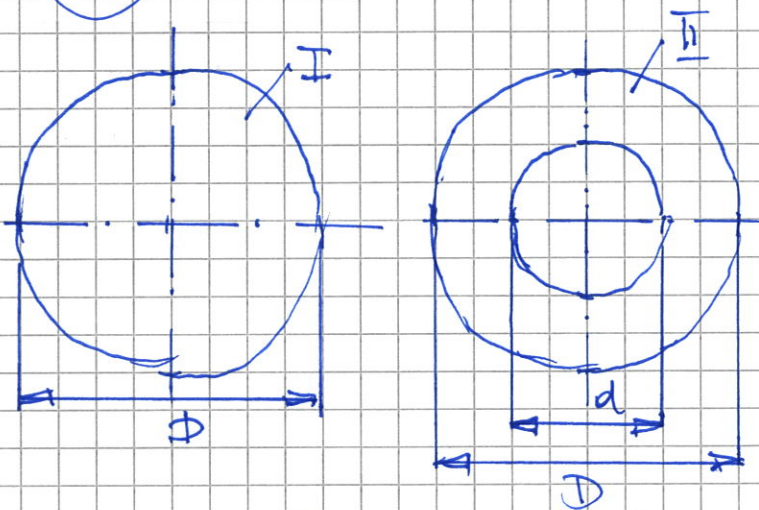
$\varphi_{max} = \frac{M \cdot l \cdot 32}{G \cdot \pi d^4}$

$$d \geq \sqrt[4]{\frac{32 M \cdot l}{\pi \cdot G \cdot \varphi_{\max}}} \quad \varphi \text{ (radians)}$$

$$d \geq \sqrt[4]{\frac{180 \cdot 32 \cdot M \cdot l}{\pi \cdot G \cdot \varphi_{\max}}} \quad \varphi \text{ (degrees)}$$

$d(\varphi_{\max}) > d(\tau_{\max})$

Ex. 2 / w. 2



$D, d = 0,5D, \gamma, l$
 $d = \alpha \cdot D, \alpha = 0,5$

~~W~~ - weight
 τ - stress

$\frac{\tau_{II}}{\tau_I} = \epsilon$ $\frac{W_{II}}{W_I} = ?$

$W_I = \frac{\pi D^2}{4} \cdot \gamma \cdot l$

$W_{II} = \frac{\pi}{4} (D^2 - d^2) \cdot \gamma \cdot l$

~~$\frac{W_{II}}{W_I} = \frac{D^2 - d^2}{D^2}$~~

$= \frac{D^2 - 0,25D^2}{D^2} = 1 - 0,25 = 0,75$ (75%)

~~$\frac{W_{II}}{W_I} = \frac{\pi D^2 \gamma \cdot l}{4} \cdot \frac{4}{\pi (D^2 - d^2) \gamma \cdot l}$~~

~~$\frac{W_{II}}{W_I} = \frac{\pi (D^2 - d^2) \gamma \cdot l}{4} \cdot \frac{4}{\pi D^2 \gamma \cdot l}$~~

$\frac{W_{II}}{W_I} = \frac{D^2 - d^2}{D^2} = \frac{D^2 - (0,5D)^2}{D^2} =$

$\tau_I = \frac{M}{W_{eI}}$

$\tau_{II} = \frac{M}{W_{eII}}$

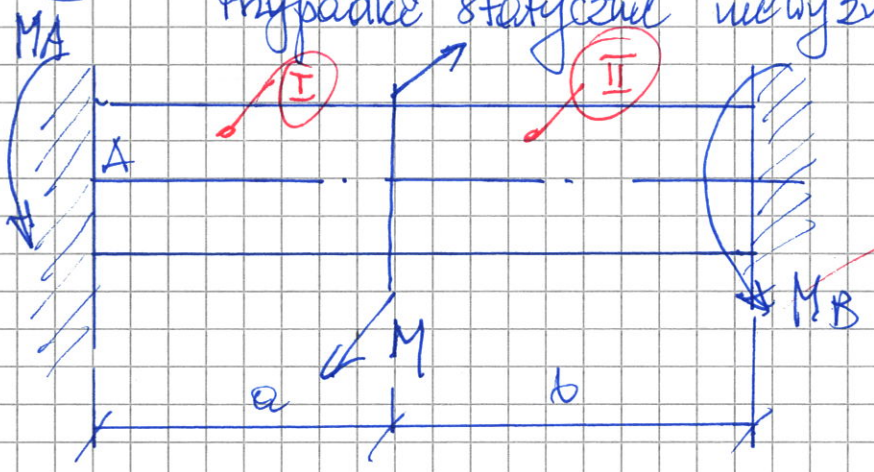
$W_{eI} = \frac{\pi D^3}{16}$

$W_{eII} = \frac{\pi}{16} (D^3 - d^3)$

$\frac{\tau_{II}}{\tau_I} = \frac{D^3 - d^3}{D^3} = \frac{D^3 - (0,5D)^3}{D^3} =$

$= 1 - 0,125 = 0,875$ (87,5%)

E Statically indeterminate problems (hyperstatic)
 Przypadki statycznie niewyznaczone



M, a, b, d, G

$T_I, T_{II} = ?$

$$I_0 = \frac{\pi d^4}{32}$$

$$K_p = \frac{\pi d^3}{16}$$

I static eqs.

$$\sum M_i = -M_A + M - M_B = 0$$

2 reactions - 1 st. eq \Rightarrow 1x hyperstatic

II Geometrical eq.

$$\phi_I + \phi_{II} = \phi = 0$$

III Physical eqs.

$$\phi_I = \frac{-M_A \cdot a}{GI_0} \quad ; \quad \phi_{II} = \frac{-(M_A + M) \cdot b}{GI_0}$$

$$-\frac{M_A \cdot a}{GI_0} + \frac{(M - M_A) \cdot b}{GI_0} = 0 \quad | \cdot GI_0$$

$$-M_A \cdot a + M \cdot b - M_A \cdot b = 0$$

$$-M_A(a+b) + M \cdot b = 0 \quad \Rightarrow \quad M_A = M \frac{b}{a+b}$$

$$T_I = \frac{M_A}{\frac{\pi d^3}{16}} = \frac{16 M_A}{\pi d^3} = \frac{16 M b}{\pi d^3 (a+b)}$$

$$T_{II} = \frac{M - M_A}{\frac{\pi d^3}{16}} = \frac{16 (M - M_A)}{\pi d^3} = \frac{16 M a}{\pi d^3 (a+b)}$$